

A580: Topics in Extragalactic Astronomy -Galaxy Disks

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National Research Council Canada Conseil national de recherches Canada



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Beauty in Rotation!

air NGC 3314









- Disks are thought to be inherently fragile. If so, how did a violent universe of hierarchical mergers produce a local galaxy population greatly dominated by disks?
- Why do galaxies rotate? How do they acquire angular momentum? Why do disks and dark matter haloes do not seem to exchange angular momentum?
- What did the first disks look like? Are they large? Are they cold? Is mass assembly efficient enough in the early Universe?
- They are beautiful!!



- Photometric properties
- Kinematics/dynamics
- Scaling Relations
- Formation Models
- Disks at high redshifts



Photometric Properties





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Exponential Law

Disk radial profile follows:

$$\Sigma(r) = \Sigma_0 \exp(-r/r_d)$$

The total disk luminosity is obtained by integrating

$$L_{tot} = 2\pi \int_0^{2\pi} \int_0^\infty \Sigma_0 \exp(-r/r_d) r dr d\theta$$

$$L_{tot} = 2\pi r_d^2 \Sigma_0$$

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Exponential Law

Disk radial profile follows:

 $\Sigma(r) = \Sigma_0 \exp(-r/r_d)$

Disk central surface brightness

The total disk luminosity is obtained by integrating

$$L_{tot} = 2\pi \int_0^{2\pi} \int_0^\infty \Sigma_0 \exp(-r/r_d) r dr d\theta$$

$$L_{tot} = 2\pi r_d^2 \Sigma_0$$

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Exponential Law

Disk radial profile follows:

$$\Sigma(r) = \Sigma_0 \exp(-r r_d) -$$

Disk scale length (often simply called "size")

The total disk luminosity is obtained by integrating

$$L_{tot} = 2\pi \int_0^{2\pi} \int_0^\infty \Sigma_0 \exp(-r/r_d) r dr d\theta$$

$$L_{tot} = 2\pi r_d^2 \Sigma_0$$

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Exponential Law

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Exponential Law

Disk radial profile follows:

$$\Sigma(r) = \Sigma_0 \exp(-r/r_d)$$

The total disk luminosity is obtained by integrating

$$L_{tot} = 2\pi \int_{0}^{2\pi} \int_{0}^{\infty} \Sigma_{0} \exp(-r/r_{d}) r dr d\theta$$
 which gives

 $L_{tot} = 2\pi r_d^2 \Sigma_0$

Infinite disks are not real of course!

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Exponential Law

Disk radial profile follows:

$$\Sigma(r) = \Sigma_0 \exp(-r/r_d)$$

The total disk luminosity is obtained by integrating

$$L_{tot} = 2\pi \int_0^{2\pi} \int_0^\infty \Sigma_0 \exp(-r/r_d) r dr d\theta$$

$$L_{tot} = 2\pi r_d^2 \Sigma_0$$



Extended UV Disks



(Thilker et al. 2007)

Fig. 4.— FUV-NIR imagery and classification contours for Type 2 XUV-disk galaxy NGC 2090. We observe a rather large blue LSB zone, which dominates the spatial extent of the galaxy despite being of low (optical) surface brightness. The image passbands and contour types are identical to those of Fig. 3. The field of view spans \$3D_ {25}=12.9^{\prime }=42.4\$ kpc at 11.3 Mpc.

Fig. 3.— FUV-NIR imagery and classification contours for XUVdisk galaxy NGC 5055 (M63), a prototype for our Type 1 class. On the left we show the GALEX FUV image of the galaxy. On the right we show the 2MASS \$K {s}\$ -band, DSS2-red, and DSS2blue imaging (as RGB channels) for an identical field of view (\$3D {25}=37.5^{\prime }=89.4\$ kpc at 8.2 Mpc). Contours are the same on both images. At the green line, the FUV surface brightness (corrected for Galactic foreground extinction and measured at 1 kpc resolution) is \$\mu _{\mathrm{F}\,\mathrm UV,=27.25\$ AB mag arcsec-2. This is the position at which (apparent?) star formation threshold mechanisms are thought to become important. The yellow contour encloses 80% of the \$K {s}\$ -band luminosity of the galaxy, defining the effective extent for the old stellar population. Note the structured UV-bright emission features beyond the green (UV) contour, which give this galaxy the Type 1 XUV-disk designation.







Figure 1. The three main disk types: Type I, Type II, and Type III (from left to right). Azimuthally averaged, radial SDSS surface brightness profiles in the g'- (triangles) and r'-band (circles) overlaid by r'-band exponential fits to the individual regions: single disk; inner and outer disk.

(Pohlen et al. 2008, ASP Conf. Series, Vol. 390, 247)



"Freeman's Law" For Galaxy Disks



B-band central surface brightness of disks is constant

μ_B(0) ~ 21.65 mag/ arcsec²

(Freeman 1970)



"Freeman's Law" For Galaxy Disks



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Millenium Galaxy Catalog Surface Brightness Distribution



Note change in µ with absolute magnitude

(Driver et al. 2005, MNRAS, 360, 81)



SDSS Disk Luminosity-Size Relation





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Low Surface Brightness Disks



(Bothun 1997, PASP, 109, 745)



Disk Size Function



$$\Phi(M, \log R_d) dM d \log R_d = \sum_{i=1}^N \frac{1}{V_{max,i}}$$

1000 Sb - Sdm galaxies

(de Jong & Lacey 2000, ApJ, 545, 781)

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Disk Size Function



115,000 disks in SDSS

(Simard 2008, ApJ, in prep.)

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Are Disks Circular?

First compute distribution of apparent axial ratios using:

$$q_{am} = \left(\frac{1-e}{1+e}\right)^{1/2}$$

$$e \equiv (e_{+}^{2} + e_{\times}^{2})^{1/2}$$



12,000 galaxies in SDSS

(Ryden 2005, ApJ, 601 214)



Are Disks Circular?

Assume disk thickness follows a Gaussian distribution:

$$f(\gamma) \propto \exp\left[-\frac{(\gamma - \mu_{\gamma})^2}{2\sigma_{\gamma}^2}\right]$$

Assume disk ellipticity follows a lognormal distribution:

$$f(\epsilon) \propto \frac{1}{\epsilon} \exp\left[-\frac{(ln\epsilon-\mu)^2}{2\sigma^2}\right]$$

Pick values for (μ_{γ} , σ_{γ} , μ and σ), compute resulting apparent axial ratio, repeat to build a distribution and compare with observed distribution. Repeat. Best values:

$$\mu_{\gamma} = 0.222 \quad \sigma_{\gamma} = 0.057 \quad \mu = -1.85 \quad \sigma = 0.89$$
$$\bullet \quad \bullet \quad e(\text{mode, median, mean}) = (0.071, 0.16, 0.21)$$



Vertical Disk Profile

Consider the luminosity density of 3D disk of the form:

$$L(R, z) = L_0 e^{-R/h_R} f(z)$$
$$f(z) = \operatorname{sech}^{2/N} (Nz/z_0)$$

(Yoachim & Dalcanton 2006, AJ, 131, 226)

where N = 1 for a self-gravitating, isothermal sheet. For $N \rightarrow \infty$,

$$f(z) \propto e^{-z/h_z}$$
 where $h_z = z_0/2$

N ≈ 2 seems to reproduce real disks (van der Kruit 1988).



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Thick Disks Form Through Accretion?





-5 0 5 10 15

R (arcsec)

Diversity of thick disk kinematics below V_c = 120 km/s

> V_c = 120 km/s

(Yoachim & Dalcanton 2006)



An interesting question (Tóth & Ostriker 1992, ApJ, 389, 5):

Given the thinness and coldness of disks, what kind of limit does this put on the current rate of infall of satellites onto spiral galaxies?

Answer: No more than 4% of the mass inside the solar radius can have been accreted within the last 5 Gyrs (expect > 28% in Ω = 1)

THIS CANNOT BE RIGHT! What is the solution? (An open universe with Λ is part of the answer ...)

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How Can Thin Disks Even Exist?





Satellites come in on radial orbits \rightarrow efficient angular momentum shedding + mass stripping \rightarrow $\Delta H \propto (M_{sat}/M_{disk})^2$

and not the TO92 regime of

 $\Delta H \propto (M_{sat}/M_{disk})$

(Hopkins et al. 2008, astro-ph/0806.2861)

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Color Distribution





B - R

Optical-nearIR colors + Bruzual-Charlot Stellar Population Models

(Bell & de Jong 2000, MNRAS, 312, 497)

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Mass-to-Light Ratios

(Bell & de Jong 2001, ApJ, 550, 212)

Salpeter IMF



(open)



Metallicity Trends





Dust in Disks



Dotted line = full sample

Solid line = B/T cut

(Driver et al. 2007, MNRAS, 379, 1022)



 $M_d(b/a) \approx M_d(1) + (2.5)(0.6)\log_{10}(b/a)$ $R_d(b/a) \approx R_d(1)/[1.0-(0.2)\log_{10}(b/a)]$

Disk (B/T < 0.3) Luminosity-Size versus Disk b/a





Disk (B/T < 0.3) Luminosity-Size versus Disk b/a

versus Disk b/a Disk (B/T < 0.3) Luminosity-Size versus Disk b/a

B/T < 0.3

Internal Absorption in Late-Type Disks

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 $M_{d}(b/a) \approx M_{d}(1) + (2.5)(0.0)\log_{10}(b/a)$ $R_{d}(b/a) \approx R_{d}(1) / [1.0 - (0.0)\log_{10}(b/a)]$

Disk (B/T > 0.5) Luminosity-Size versus Disk b/a





B/T > 0.5

Disk (B/T > 0.5) Luminosity-Size versus Disk b/a

Disk (B/T > 0.5) Luminosity-Size versus Disk b/a

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Internal Absorption in Early-Type Disks



Dust in Disks

17879 SDSS galaxies, 0 <= b/a <= 0.25



74500 SDSS galaxies, 0.50 < b/a <= 0.75



(30) (3)

38963 SDSS galaxies, 0.25 < b/a <= 0.50

88365 SDSS galaxies, 0.75 < b/a <= 1.00



(Simard 2008, in prep.)



Bars in Disks



BAR FRACTION IN THE ${\cal H}$ BAND and Optical Catalogs

Bar Class (1)	H BAND		RC3		CAG	
	Fraction (%) (2)	Number (3)	Fraction (%) (4)	Number (5)	Fraction (%) (6)	Number (7)
SB	56	105	35	65	27	44
SAB	16	30	30	56	4	6
SB+SAB	73	135	65	121	30	50
SA	27	51	35	65	70	116
SA+SAB	44	81	66	121	74	122

Eskridge et al. 2000, AJ, 119, 536

Cold disks are very unstable

Disks with large radial velocity dispersion are immune to bars

Bars are transient phenomena **NRC** · **CNRC** Herzberg Institute of Astrophysics

Bar Dissolution and Bulge Formation



(Norman et al. 2006, ApJ, 462, 114)

Gas is funneled to the center of the disk, it triggers a starburst, forms a bulge. Once bulge is formed, stars on radial orbits in the bar scatter off and bar dissolves.



Warps in Disks



Figure 2. (a) The warp at t = 400 in our simulation; its morphology closely resembles the observed H_I warp of NGC 4013 (reproduced with permission Bottema 1996) shown in (b). The length unit shown is the scalelength R_d of the exponential disc. Note that, we have oriented the model so that the inner $(R < 3R_d)$ disc lies in the x-y plane, which is perpendicular to the paper.

Evolving halo due to cosmic infall creates different torques on inner and outer disks

(Shen & Sellwood 2006, MNRAS, 370, 2) NRC CNRC

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Age of MW Disk in Cosmological Context



Age determination based on white dwarf cooling sequence

(Hansen et al. 2007, ApJ, 671, 380)



Kinematics / Dynamics

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Internal Kinematics -Observations

- Challenges
 - Sensitivity
 - Spatial coverage
 - Spatial resolution ("beamwidth")
 - Spectral resolution
- Optical
 - Slitless
 - Slit Spectroscopy
 - Long (minor/major axis, drift-scan)
 - Multi
 - Integral Field Spectroscopy
 - Fabry-Perot
 - Fiber bundles
 - Image slicers
- Radio
 - Single dish
 - Interferometry



Internal Kinematics -Observations



Example of an observational challenge :

The subtle effects of beamwidth smearing

Sofue & Rubin 2001, ARAA, 39, 137

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Internal Kinematics -Tracers

Optical

- Emission-lines (gaseous)
 - [OII] 3727,3729
 - Hb 4861
 - [OIII] 4959,5007
 - Ha 6562
 - NII
- Absorption-line (stellar)
 - Mg triplet 5167,5173,5184
 - G-band 4300
 - Ca triplet 8498, 8542, 8662
 - Ca H + K 3934, 3968
- Planetary nebulae ([OIII] 5007)

Radio

- Neutral hydrogen (HI)
- Carbon Monoxide (CO)
- SiO, OH and H₂O Masers in circumstellar envelopes



Internal Kinematics -Tracers



Hα emission line (high spatial resolution)

GHASP Survey

Garrido et al. 2002, A+A, 387, 821



Internal Kinematics -Tracers



Ca H and K lines (Stellar velocity dispersion)

Kobulnicky & Gebhardt 2002, AJ, 119, 1608

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Internal Kinematics -Tracers



Neutral Hydrogen (HI)

Kobulnicky & Gebhardt 2002, AJ, 119, 1608



Internal Kinematics -Tracers

Carbon Monoxide (CO)

"arctan velocity field"

(Think ALMA ...)

Sofue et al. 1999, ApJ, 523, 136



Internal Kinematics -Tracer Comparison





Internal Kinematics -Tracer Comparison



[OII] versus HI (very different spatial resolutions)

Kobulnicky & Gebhardt 2002, AJ, 119, 1608

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Internal Kinematics -Types of Motion

- Basic components ("Easier to measure")
 - Rotation ($V_c \text{ or } V_{rot}$)
 - Dispersion (σ)

• Peculiar motions ("Hard to measure")

- Massive black holes and circumnuclear rotation
- Counter-rotation
 - Nuclear disks
 - Extended disks
 - Decoupled cores
- Resonance rings
- Lopsided gas
- Bars, warps, supershells
- Interactions and mergers

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Internal Kinematics -Rotation Curves



The outer parts of disk galaxy rotation curves are "flat", i.e., M(r) α r

One of the best pieces of evidence for the existence of dark matter

Kent 1987, AJ, 93, 816 RC-CRC Herzberg Institute of Astrophysics

Mass-Halo Degeneracy





Can rotation curves help us determine the shape of dark matter halos?

$$\rho(r) = \frac{\rho_0}{[c + (r/r_0)^{\gamma}][1 + (r/r_0)^{\alpha}]^{(\beta - \gamma)/\alpha}},$$

where ρ_0 and r_0 are characteristic density and radius. The constant c forces the presence of a flat core. The parameters, α , β , and γ , determine the shape of the halo profile.

Four different models for (c, α , β , γ):

- 1. Pseudo-isothermal sphere (1, α != 0, 2, 2)
- 2. Navarro-Frenk-White (0, 1, 3, 1)
- 3. Burkert (1, 2, 3,1)
- 4. Klypin et al. (0, 2, 3, 0.2)



Rotation Curves -Cores





460 km/s

300 km/s

NGC 3109

Rotation Curves -Cores

<u>Very</u> inhomogenous emission-line gas distribution

Slit observations are not adequate

Use Integral Field Spectroscopy

Blais-Ouellette et al. 2001, AJ, 121, 1952

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Rotation Curves -Cores



Blais-Ouellette et al. 2001, AJ, 121, 1952

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(km/s)



Figure 3. Inner logarithmic slope (γ) of the CDM halos. Note that N-body simulations give $\gamma \geq 1$.

Blais-Ouellette et al. 2002, astro-ph/0203146

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Velocity Fields - Cores





More mass in the core of NFW profiles than allowed in full 2D IFU velocity fields

Kuzio de Naray et al. 2008, ApJ, 676, 920

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Scaling Relations

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Scaling Relations

- Galaxies have a fundamental set of physical properties
 - Mass
 - Size
 - Angular momentum
- These physical properties translate into observables
 - Luminosity or mass-to-light ratio
 - Velocity (rotation and/or dispersion)
 - Size
 - Half-light radius (R_{hl} or R_e)
 - Disk scale length (R_d or h)
- <u>Basic questions</u> : What are the underlying relations (if any) among galaxy properties and what do they tell us about galaxy formation?

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Scaling Relations -Observables

- Luminosity
 - Integrated
 - Disk-only (model dependent)
 - Choice of bandpass
- Velocity
 - Integrated?
 - Rotation curve: where do you measure V_c ?
 - Peak
 - Flat region
 - Intermediate radius (~ 2.2 R_d)
- Size
 - Half-light radius
 - Disk scale length (model dependent)



- Opik (1922) : Assume all spiral nebulae have same M/L => distance to M31 of 450 kpc (!!!)
 - (We know V² = GM/R. Assume L ~ R², then L \propto V⁴)
- Hubble (1936) : Failed to recognize use as a distance indicator
- Roberts (1962)
- Balkowski et al. (1974) : HI width is correlated with luminosity
- Tully & Fisher (1977): Hey! <u>Please</u> use correlation as a distance indicator!!! ⁵⁹



Tully-Fisher Relation -History



Tully & Fisher 1977, A&A, 54, 66



Tully-Fisher Relation -Modern Version



Simard & Pritchet 1998, ApJ, 505, 96

Using data for 1355 galaxies from Mathewson et al. 1992, ApJS, 81, 413

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Tully-Fisher Relation -Biases and Corrections

- Inclination measurements from direct optical images
 - Disks may not be circular ($q_0 = 0.11 0.20$)
 - **Disks may not have constant ellipticity**
 - **Disks are not infinitely thin (+seeing)**

$$i = \cos^{-1} \sqrt{\frac{(b/a)^2 - q_0^2}{1 - q_0^2}},$$

• Spectral resolution : $W_{raw}^2 = W_{gal}^2 + W_{res}^2$

• Internal extinction
$$A_i^{\lambda} = \gamma_{\lambda} \log \left(\frac{a}{b} \right)$$

$$A_i^{\lambda} = \gamma_{\lambda} \log\left(\frac{a}{b}\right) \,,$$

 $W_R^2 = W_{20}^2 + W_t^2 - 2W_t W_{20} \left(1 - e^{-W_{20}^2/W_c^2}\right)$ Turbulence (HI)

$$-2W_t^2 e^{-W_{20}^2/W_c^2} + 4W_{dwarf}^2$$
,

Malmquist bias

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Tully-Fisher Relation -Slope



Slope steepens towards redder bandpasses (real?)

Jacoby et al. 1992, PASP, 104, 599



Tully-Fisher Relation -Slope



"Pruning"

Verheijen 2001 ApJ, 563, 694



Tully-Fisher Relation -Scatter



Which radius should be used for V?

Verheijen 2001 ApJ, 563, 694
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Tully-Fisher Relation -Scatter



Samples for TF distances are "pruned"

Choice of internal extinction and fitting method can change results dramatically

Kannappan et al. 2002, AJ, 123, 2358 ₆₆





Sa galaxies exhibit offsets:

> 0.76 mag in R 0.95 mag in B 1.20 mag in U

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Tully-Fisher Relation -Scatter



Pruning by morphology and kinematics





Scatter depends on morphology





Systematic dependence of residuals on color and EW(Hα)





Color correction reduces scatter





Including all types and luminosities increases scatter

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Tully-Fisher Relation -Scatter



Scatter is neither correlated with surface brightness nor with gas consumption



Tully-Fisher Relation -Turbulence



Turbulence introduces an offset

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Tully-Fisher Relation -Barred Galaxies



Barred galaxies follow same TFR as "normal" galaxies

Federspiel 1999 Ph.D. Thesis



Tully-Fisher Relation -Galaxy Pairs



Galaxy pairs follow same TFR as "normal" galaxies except for a few outliers

Barton et al. 2001, AJ, 121, 625



Tully-Fisher Relation -Surface Brightness



LSB galaxies follow same TFR as HSB galaxies -

 $(M/L)_{LSB} \sim 2 (M/L)_{HSB}$

Since $M \sim V^2h$, then $h_{LSB} \sim 2h_{HSB}$

Zwaan et al. 1995 MNRAS, 273, L35

Figure 2. Tully–Fisher relations binned according to morphological type. The solid circles are LSB galaxies, and the open circles are HSB galaxies from Broeils. The line is a fit to all LSB galaxies.

Herzberg Institute of Astrophysics Tully-Fisher Relation -Summary

- Slope steepens towards redder bandpasses (internal extinction?)
- Sa galaxies exhibit a systematic offset
- TFR scatter depends on:
 - Radius at which V_c is measured
 - Bandpass (internal extinction)
 - Morphology
 - Galaxy color
 - Star formation rate
- LSB galaxies follow the same relation as HSB galaxies
- Galaxies in pairs follow the same relation as "normal" galaxies
- Barred galaxies follow same relation as "normal galaxies"



Tully-Fisher Relation -A Distance Indicator



Final results of HST H₀ Key Project

TFR is a <u>secondary</u> distance indicator

Freedman et al. 2001, ApJ, 553, 47



Formation Models





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N-Body Simulation circa 2000



z: 49.5

Ζ

X

X

Stars color-coded according to age with blue (young) and red (old)

Navarro & Steinmetz 2000, ApJ, 538, 477



For a singular isothermal sphere the density profile is just

$$\rho(r) = \frac{V_{\rm c}^2}{4\pi G r^2},\tag{1}$$

$$r_{200} = \frac{V_{\rm c}}{10H(z)}; \quad M = \frac{V_{\rm c}^2 r_{200}}{G} = \frac{V_{\rm c}^3}{10GH(z)},$$
 (2)

We assume that the mass which settles into the disc is a fixed fraction m_d of the halo mass. The disc mass is then

$$M_{\rm d} = \frac{m_{\rm d} V_{\rm c}^3}{10 GH(z)} \approx 1.7 \times 10^{11} h^{-1} \,\mathrm{M}_{\odot} \left(\frac{m_{\rm d}}{0.05}\right) \left(\frac{V_{\rm c}}{250 \,\mathrm{km \, s^{-1}}}\right)^3 \left[\frac{H(z)}{H_0}\right]^{-1}.$$
(4)

Fall & Efstathiou 1980, MNRAS, 193, 189; Mo et al. 1998, MNRAS, 295, 319



$$\Sigma(R) = \Sigma_0 \exp(-R/R_d).$$
⁽⁵⁾

Here R_d and Σ_0 are the disc scalelength and central surface density, and are related to the disc mass through

$$M_{\rm d} = 2\pi\Sigma_0 R_{\rm d}^2. \tag{6}$$

If the gravitational effect of the disc is neglected, its rotation curve is flat at the level V_c and its angular momentum is just

$$J_{\rm d} = 2\pi \int V_{\rm c} \Sigma(R) R^2 dR = 4\pi \Sigma_0 V_{\rm c} R_{\rm d}^3 = 2M_{\rm d} R_{\rm d} V_{\rm c}.$$
 (7)

We assume this angular momentum to be a fraction j_d of that of the halo, i.e.

$$J_{\rm d} = j_{\rm d} J,\tag{8}$$



and we relate J to the spin parameter λ of the halo through the definition

$$\lambda = J|E|^{1/2}G^{-1}M^{-5/2},\tag{9}$$

where *E* is the total energy of the halo. Equations (7) and (8) then imply that

$$R_{\rm d} = \frac{\lambda G M^{3/2}}{2V_{\rm c} |E|^{1/2}} \left(\frac{j_{\rm d}}{m_{\rm d}}\right). \tag{10}$$

The total energy of a truncated singular isothermal sphere is easily obtained from the virial theorem by assuming all particles to be on circular orbits:

$$E = -\frac{GM^2}{2r_{200}} = -\frac{MV_c^2}{2}.$$
(11)



Theoretical Disk Formation (cont'd)

$$R_{\rm d} = \frac{1}{\sqrt{2}} \left(\frac{j_{\rm d}}{m_{\rm d}} \right) \lambda r_{200}$$

$$\approx 8.8h^{-1}\,\mathrm{kpc}\left(\frac{\lambda}{0.05}\right)\left(\frac{V_{\mathrm{c}}}{250\,\mathrm{km\,s^{-1}}}\right)\left[\frac{H}{H_0}\right]^{-1}\left(\frac{j_{\mathrm{d}}}{m_{\mathrm{d}}}\right),\qquad(12)$$

$$P(\lambda)d\lambda = \frac{1}{\sqrt{2\pi\sigma_{\lambda}}} \exp\left[-\frac{\ln^2\left(\lambda/\lambda_{\rm med}\right)}{2\sigma_{\lambda}^2}\right] \frac{d\lambda}{\lambda}.$$
 (5)

The median $\lambda_{\rm med}$ and dispersion (in $\ln \lambda$) σ_{λ} are found to depend remarkably weakly on the cosmology, halo mass, or initial spectrum of density fluctuations (e.g., Barnes & Efstathiou 1987; Warren et al. 1992; Cole & Lacey 1996), with typical values $\lambda_{\rm med} \approx 0.04$ and $\sigma_{\lambda} \approx 0.5-0.6$.

σ(In λ)_{theo} ≈ 1.5x σ(In λ)_{obs}



$$R_{\rm d} = \frac{1}{\sqrt{2}} \left(\frac{j_{\rm d}}{m_{\rm d}} \right) \lambda r_{200}$$

$$\approx 8.8 h^{-1} \, \rm kpc \left(\frac{\lambda}{0.05} \right) \left(\frac{V_{\rm c}}{250 \,\rm km \, s^{-1}} \right) \left[\frac{H}{H_0} \right]^{-1} \left(\frac{j_{\rm d}}{m_{\rm d}} \right), \quad (12)$$

$$P(\lambda)d\lambda = \frac{1}{\sqrt{2\pi\sigma_{\lambda}}} \exp\left[-\frac{\ln^2\left(\lambda/\lambda_{\rm med}\right)}{2\sigma_{\lambda}^2}\right] \frac{d\lambda}{\lambda}.$$
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σ(In λ)_{theo} ≈ 1.5x σ(In λ)_{obs}

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Observed Width of Disk Size Distribution

M _{B,disk}	Ν	<log<sub>10 R_d></log<sub>	$\sigma(\log_{10} R_d)$
[-22.0,-21.5]	6246	0.82	0.12
[-21.5,-21.0]	12497	0.73	0.12
[-21.0,-20.5]	15061	0.64	0.12
[-20.5,-20.0]	13186	0.56	0.12
[-20.0,-19.5]	9593	0.48	0.13
[-19.5,-19.0]	6212	0.40	0.15
[-19.0,-18.5]	3588	0.36	0.16
[-18.5,-18.0]	2132	0.30	0.18

SDSS Disk subsample: B/T ≤ 0.3 b/a ≥ 0.5

Predicted σ(log₁₀ R_d) ≈ 0.20-0.25

Feedback at high λ ? Secular instability at low λ ?

(Simard 2008, in prep.)



Tully-Fisher Relation - Origin

$$M_{\rm d} = \frac{m_{\rm d} V_{\rm c}^3}{10GH(z)} \approx 1.7 \times 10^{11} h^{-1} \,{\rm M}_{\odot} \left(\frac{m_{\rm d}}{0.05}\right) \left(\frac{V_{\rm c}}{250 \,{\rm km \, s^{-1}}}\right)^3 \left[\frac{H(z)}{H_0}\right]^{-1} \text{ et al. 1998}$$

$$MIRAS, 295, 319$$

$$M_{\rm d} \equiv M_{\rm d}/L_{\rm d} \text{ (in solar units)}$$

$$L_{\rm d} = A \left(\frac{V_{\rm c}}{250 \,{\rm km \, s^{-1}}}\right)^{\alpha},$$
where $\alpha = 3$ is the slope, and
$$A = 1.7 \times 10^{11} h^{-1} {\rm L}_{\odot} \Upsilon_{\rm d}^{-1} \left(\frac{m_{\rm d}}{0.05}\right) \left[\frac{H(z)}{H_0}\right]^{-1}$$
88

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Tully-Fisher Relation -Origin





A Disk Stability Criterion

$$\epsilon \equiv \frac{V_{max}}{(GM_d/R_d)^{1/2}} \le 1.1$$



If self-gravity of disk dominates rotational support, then disk will be unstable leading to the formation of a bar or a bulge

$$\lambda' \equiv \lambda j_d / m_d$$

(Efstathiou et al. 1982, MNRAS, 199, 1069)



``Angular Momentum Catastrophe"



Simulations reproduce slope but not zeropoint

Navarro & Steinmetz 2000, ApJ, 538, 477



``Angular Momentum Catastrophe"



Specific angular momenta of simulated disks are too low

Navarro & Steinmetz 2000, ApJ, 538, 477

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Angular Momentum Catastrophe -Resolved?

Decreasing particle resolution in isolated CDM halo



Disk-halo transfer due to halo clumpiness?



(Mayer et al. 2008, arXiv: 0801.3845; Governato et al. 2008, arXiv: 0801.1707)



orbits

Disks in Mergers



<u>Gas-rich</u> mergers actually produce disks!

(Robertson et al. 2006, ApJ, 645, 986)



Disks in Mergers



Significant disk component in mergers (f_{disk} = 0.76)

(Hopkins et al. 2008, arXiv: 0806.1739)



Origin of Exponential Disks



Angular momentum transport?

Initially,

j_{gas} ∝ r^{1.0} j_{DM}~ 0

(Mayer et al. 2007, MNRAS, 375, 53)

Resulting surface density of gas disks is not exponential ...



Origin of Exponential Disks

$$\frac{\partial \Sigma_g}{\partial t} = -\frac{1}{r} \frac{\partial}{\partial r} \left[\frac{(\partial/\partial r)(\nu \Sigma_g r^3 d\Omega/dr)}{d/dr(r^2 \Omega)} \right] - \psi_*$$



Viscous evolution?

Caused by noncircular motions and turbulence in differentially rotating disks - no shear = no viscous evolution

(Lin & Pringle 1987, ApJ, 320, L87; Bell 2002, ApJ, 581, 1013



Disks at High Redshifts

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Luminosity-Size Relation at z < 1





-16

-18 -16



No surface brightness evolution?



(Simard et al. 1999, ApJ, 519, **563)**



Luminosity-Size Relation at z < 1



Selection effects are very significant

1 mag
brighter in Vband surface
brightness at z
1

(Barden et al. 2005, ApJ, 635, 959)


Disk Size Function at z < 1



Size function remains constant out to z ~ 1 if selection window is shifted by ~1 mag with z.

(Sargent et al. 2007, ApJ, 172, 434)



Disk Size Function at z < 1

Are selection effects "Lagrangian" or "Eulerian"?



(Kanwar et al. 2008, ApJ, 682, 907)



Tully-Fisher Relation at z < 1

IMAGES Survey

2D velocity fields with VLT/GIRAFFE



Blue = Rotating disks Green = Perturbed rotators Red = "Kinematically complex" (Puech et al. 2008, A&A, 484, 173)

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Tully-Fisher Relation at z < 1



DEEP2 Survey

Multi-slits with Keck/DEIMOS

Galaxies are "settling down" onto a fundamental TF relation

(Kassin et al. 2007, ApJ, 660, L35)



Very Large Disks



HDF-South FIRES

z = 2 - 3

(Labbé et al. 2003, 591, L95)



Very Large Disks

 TABLE 1

 PROPERTIES OF HIGH-REDSHIFT DISK GALAXIES IN THE HDF-S

Galaxy (1)	$K_{s, ext{tot}}$ (2)	z (3)	$M_{B, \mathrm{rest}}$ (4)	$\mu_{0,B,\mathrm{rest}}$ (5)	(arcsec) (6)	(arcsec) (7)	(arcsec) (8)	е (9)
302	19.70	1.439ª	-22.70	19.70	0.89	0.70	0.86	0.46
267	19.98	1.82	-22.88	19.92	0.75	0.74	0.88	0.37
257	20.25	2.027^{a}	-23.08	19.53	0.74	0.74	0.84	0.36
657	20.68	2.793ª	-23.56	19.33	0.76	0.70	0.74	0.18
611	20.53	2.94	-23.59	18.51	0.65 ^b	0.52	0.97	0.27
494	21.14	3.00	-23.31	18.84	0.75 ^b	0.56	0.86	0.47

HDF-South FIRES

z = 2 - 3 (7.83 kpc per arcsec)

NOTE. — Col. (1): Catalog identification numbers (see Labbé et al. 2003). Col. (2): K_s -band total magnitudes. Col. (3): Redshift. Col. (4): Rest-frame absolute *B*-band magnitudes. Col. (5): Face-on rest-frame *B*-band central surface brightnesses. Col. (6): Face-on best-fit effective radii. Col. (7): K_s half-light radii. Col. (8): I_{814} half-light radii, PSF-matched to K_s . Col. (9): Ellipticity.

^a Spectroscopic redshifts.

^b Two-component models (point+exponential).

- Sizes comparable to Milky Way
- Large Stellar Masses
- Regular K morphologies, knotty in V
- Constitute half of the most rest-frame luminous galaxies
- Number density is at least a factor of two above model predictions

(Labbé et al. 2003, ApJ, 591, L95)

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Rotation or Mergers?



"SINS" Survey

Forster-Schreiber et al. 2006, ApJ, 645, 1062



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Schmidt-Kennicutt Law at z ~ 2



Same as the local relation (!)

(Bouché et al. 2007, ApJ, 671, 303)



The End!